

Classification of the possible symmetries of the Jahn—Teller systems

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On the basis of group-theoretical approach we elaborated the complete classification of possible symmetries of extremal points of potential energy surfaces of Jahn—Teller systems for all symmetry point groups. The method originates from the principle of step-by-step splitting of degenerate irreducible representations of corresponding electron terms due to the lowering of the symmetry of the system. It was shown that the existence of some deformations cannot be the result of the Jahn—Teller effect contrary to the fact that also such an explanation is at present used.

На основе группово-теоретического подхода мы разработали полную классификацию возможных симметрий экстремальных точек поверхностей потенциальной энергии систем Яна—Теллера для всех точечных групп симметрии. Этот метод проистекает из принципа поэтапного расщепления вырожденных неприводимых представлений соответствующих электронных термов вследствие понижения симметрии системы. Было показано, что существование некоторых искажений не может являться следствием эффекта Яна—Теллера, вопреки тому, что подобное объяснение также используется в настоящее время.

The problem of symmetry of systems which originate through Jahn—Teller deformation from the original highly symmetrical system, was mostly solved by analytical methods of searching for extrema of adiabatic potential surface [1—15]. Such an attempt to solve this problem has two major disadvantages:

i) Analytical forms of adiabatic potential surfaces are determined by perturbation theory, using the Taylor expansion of vibronic Hamiltonian. The number of extrema from such adiabatic potential surfaces is then dependent on the order of perturbation theory used as well as on the number of terms in the Taylor expansion of vibronic Hamiltonian.

ii) The problem of obtaining the extremal points from adiabatic potential of more complex forms is transformed to the problem of solving the system of nonlinear equations. We note that there does not exist any algorithm which guaranties obtaining all the extremal points.

Some authors [16—19] used the group-theoretical approach to solve this problem using the symmetry of the Jahn—Teller active coordinates. Even this approach did not give the complete solution of the problem, because the obtained results satisfy only the first order of perturbation theory.

In our work we present the complete group-theoretical analysis of symmetries of the Jahn—Teller systems of all molecular symmetry point groups.

Method and results

The Jahn—Teller effect may appear only in the nonlinear systems with such a symmetry group which possesses at least one multi-dimensional irreducible representation. Mostly this is two-dimensional degenerate *E*-representation and/or three-times degenerate *T*-representation. In the case of symmetry groups I_h and *I* yet four-times degenerate *G*-representation and five-times degenerate *H*-representation appear. Such a multi-dimensional representation splits due to the Jahn—Teller effect into the representation of lower order (until the nondegenerate representation is achieved) lowering the symmetry of the system. Therefore the problem to determine the extremal points of the Jahn—Teller systems is thus reduced to find all the symmetry point groups which satisfy the following conditions:

i) They are the subgroups of original symmetry group of the system in which the Jahn—Teller effect appears.

ii) By splitting the degenerate irreducible representations there appears at least one nondegenerate irreducible representation.

iii) The subgroups are not permitted if there exists the transition group (such that it causes the step-by-step lowering of symmetry of original group to the final subgroup) while the degeneracy is removed already in the transition group.

For our purposes still further division of such subgroups into some levels will be useful. In the first level there are groups which are the subgroups of the former original symmetry group of studied system and are not the subgroups of any other subgroups. In the second level there will be groups which are the subgroups of the original group as well as of the groups from the first level. In the third level there will be groups which are the subgroups of the original group as well as of the groups from the first and the second level, *etc.* The order of the original symmetry group is the integer multiple of the orders of its subgroups [20, 21]. This gives us the number of equivalent geometries of extremal points of adiabatic potential surface, which correspond to the system with the symmetry of the given subgroup.

The hierarchy of molecular symmetry point groups as well as their subgroups is presented in Fig. 1. The common dihedral and cyclic groups have been arranged in the figure in descending sequence from groups of the highest to the lowest order. This enables all the subgroups of a particular point group to be noted at a glance. The quotient of the order of given group and its subgroups must be a small integer greater than unity. The icosahedral

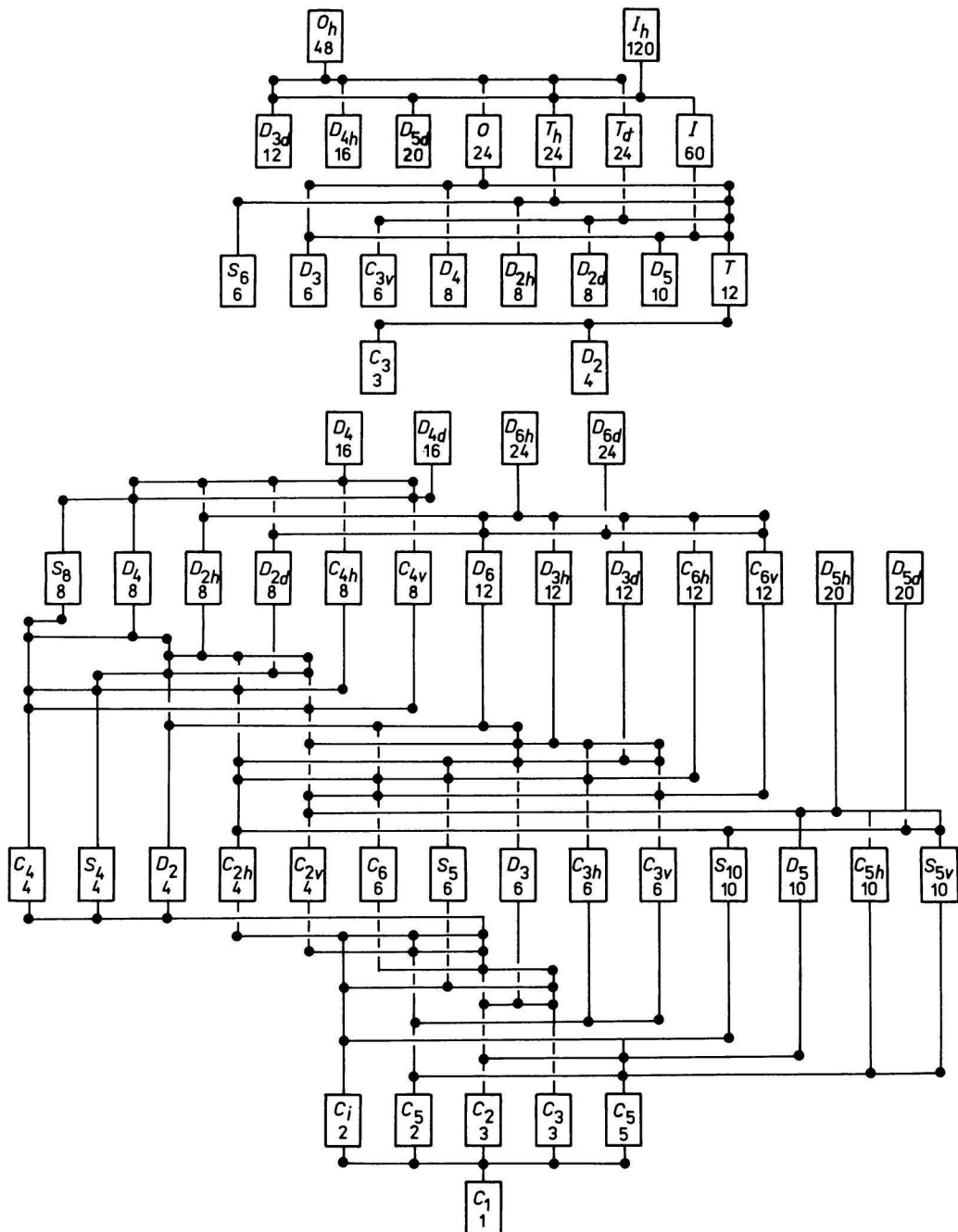


Fig. 1. Hierarchy of molecular symmetry point groups.

and cubic groups have been arranged in a similar figure and the descent in symmetry continued until correspondence with the other part of figure was achieved.

The results of group-theoretical analysis of all the symmetry point groups (containing the degenerate irreducible representations) of molecules are presented in Table 1. Subgroups fulfilling the three mentioned conditions are ordered into individual levels. For each subgroup we also present the irreducible representations to which the stable electron term can be ascribed.

Table 1

Possible symmetries of the extremal points of Jahn—Teller systems^a

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term	
O_h	48	E_g	1	D_{4h}	16	3	A_{1g}, B_{1g}	
				D_4	8	6	A_1, B_1	
			D_{2h}	8	6	A_g		
			D_{2d}	8	6	A_1, B_1		
			3	C_{2h}	4	12	A_g, B_g	
				D_2	4	12	A	
				C_2	2	24	A, B	
				C_i	2	24	A_g	
				C_s	2	24	A', A''	
				C_1	1	24	A', A''	
			E_u	1	D_{4h}	16	3	A_{1u}, B_{1u}
					D_4	8	6	A_1, B_1
		D_{2h}		8	6	A_u		
		D_{2d}		8	6	A_1, B_1		
		3		C_{2h}	4	12	A_u, B_u	
				D_2	4	12	A	
				C_2	2	24	A, B	
				C_i	2	24	A_u	
				C_s	2	24	A', A''	
				C_1	1	24	A', A''	
		T_{1g}		1	D_{4h}	16	3	A_{2g}
					D_{3d}	12	4	A_{2g}
			2	D_4	8	6	A_2	
				D_3	6	8	A_2	
D_{2d}	8			6	A_2			
D_{2h}	8			6	B_{1g}, B_{2g}, B_{3g}			
S_6	6			8	A_g			
C_{3v}	6			8	A_2			
C_{2h}	4		12	A_g, B_g				

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term
			3	D_2	4	12	B_1, B_2, B_3
				C_3	3	16	A
				C_{2v}	4	12	B_1, B_2
				C_{2h}	4	12	B_g
				C_2	2	24	A, B
				C_i	2	24	A_g
				C_s	2	24	A', A''
		T_{2g}	4	C_2	2	24	B
			1	D_{4h}	16	3	B_{2g}
				D_{3d}	12	4	A_{1g}
			2	D_4	8	6	B_2
				D_3	6	8	A_1
				D_{2d}	8	6	B_2
				D_{2h}	8	6	B_{1g}, B_{2g}, B_{3g}
				S_6	6	8	A_g
				C_{3v}	6	8	A_1
				C_{2h}	4	12	A_g, B_g
			3	D_2	4	12	B_1, B_2, B_3
				C_3	3	16	A
				C_{2v}	4	12	B_1, B_2
				C_{2h}	4	12	B_g
				C_2	2	24	A, B
				C_i	2	24	A_g
				C_s	2	24	A', A''
		T_{1u}	4	C_2	2	24	B
			1	D_{4h}	16	3	A_{2u}
				D_{3d}	12	4	A_{2u}
			2	D_4	8	6	A_2
				D_3	6	8	A_2
				D_{2d}	8	6	B_2
				D_{2h}	8	6	B_{1u}, B_{2u}, B_{3u}
				S_6	6	8	A_u
				C_{3v}	6	8	A_1
				C_{2h}	4	12	A_u, B_u
			3	D_2	4	12	B_1, B_2, B_3
				C_3	3	16	A
				C_{2v}	4	12	B_1, B_2
				C_{2h}	4	12	B_u
				C_2	2	24	A, B
				C_i	2	24	A_u
				C_s	2	24	A', A''

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term	
I_h	120	T_{2u}	4	C_2	2	24	B	
			1	D_{4h}	16	3	B_{2u}	
				D_{3d}	12	4	A_{1u}	
			2	D_4	8	6	B_2	
				D_3	6	8	A_1	
				D_{2d}	8	6	B_2	
				D_{2h}	8	6	B_{1u}, B_{2u}, B_{3u}	
				S_6	6	8	A_u	
				C_{3v}	6	8	A_2	
				C_{2h}	4	12	A_u, B_u	
			3	D_2	4	12	B_1, B_2, B_3	
				C_3	3	16	A	
				C_{2v}	4	12	B_1, B_2	
				C_{2h}	4	12	B_u	
				C_2	2	24	A, B	
				C_i	2	24	A_u	
				C_s	2	24	A', A''	
			4	C_2	2	24	B	
		T_{1g}	1	D_{5d}	20	6	A_{2g}	
				D_{3d}	12	10	A_{2g}	
			2	D_5	10	12	A_2	
				D_3	6	20	A_2	
				D_{2h}	8	15	B_{1g}, B_{2g}, B_{3g}	
				S_6	6	20	A_g	
				C_{2h}	4	30	A_g, B_g	
			3	D_2	4	30	B_1, B_2, B_3	
				C_3	3	40	A	
				C_2	2	60	A, B	
				C_i	2	60	A_g	
				C_s	2	60	A', A''	
			T_{1u}	1	D_{5d}	20	6	A_{2u}
					D_{3d}	12	10	A_{2u}
				2	D_5	10	12	A_2
					D_3	6	20	A_2
					D_{2h}	8	15	B_{1u}, B_{2u}, B_{3u}
					S_6	6	20	A_u
	C_{2h}	4		30	A_u, B_u			
3	D_2	3		40	A			
	C_2	2		60	A, B			
	C_i	2		60	A_u			
	C_s	2		60	A', A''			

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term					
G_g	1	G_g	1	T_h	24	5	A_g					
				D_{3d}	12	10	A_{1g}, A_2					
				T	12	10	A					
				D_3	6	20	A_1, A_2					
				D_{2h}	8	15	B_{1g}, B_{2g}, B_{3g}					
				S_6	6	20	A_g					
	2	G_g	2	C_{2h}	4	30	A_g, B_g					
				D_2	4	30	B_1, B_2, B_3					
				C_3	3	40	A					
				C_2	2	60	A, B					
				C_i	2	60	A_g					
				C_s	2	60	A', A''					
	G_u	1	G_u	1	T_h	24	5	A_u				
					D_{3d}	12	10	A_{1u}, A_{2u}				
		2		G_u	2	T	12	10	A			
						D_3	6	20	A_1, A_2			
						D_{2h}	8	15	B_{1u}, B_{2u}, B_{3u}			
						S_6	6	20	A_u			
						C_{2h}	4	30	A_u, B_u			
						D_2	4	30	B_1, B_2, B_3			
3		G_u		3	C_3	3	40	A				
					C_2	2	60	A, B				
					C_i	2	60	A_u				
					C_s	2	60	A', A''				
					H_g	1	H_g	1	D_{5d}	20	6	A_{1g}
									D_{3d}	12	10	A_{1g}
2	H_g	2	D_5	10	12	A_1						
			D_3	6	20	A_1						
			D_{2h}	8	15	$A_g, B_{1g}, B_{2g}, B_{3g}$						
			S_6	6	20	A_g						
			C_{2h}	4	30	A_g, B_g						
			D_2	4	30	A, B_1, B_2, B_3						
3	H_g	3	C_3	3	40	A						
			C_2	2	60	A, B						
			C_i	2	60	A_g						
			C_s	2	60	A', A''						
			H_u	1	H_u	1	D_{5d}	20	6	A_{1u}		
							D_{3d}	12	10	A_{1u}		
2	H_u	2	D_5	10	12	A_1						
			D_3	6	20	A_1						
			D_{2h}	8	15	$A_u, B_{1u}, B_{2u}, B_{3u}$						
			S_6	6	20	A_u						

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term		
T_d	24	E	3	C_{2h}	4	30	A_u, B_u		
				D_2	4	30	A, B_1, B_2, B_3		
				C_3	3	40	A		
				C_2	2	60	A, B		
				C_i	2	60	A_u		
				C_s	2	60	A', A''		
			1	D_{2d}	8	3	A_1, B_1		
				D_2	4	6	A		
				C_3	2	12	A', A''		
				T ₁	1	D_{2d}	8	3	A_2
						C_{3v}	6	4	A_2
						D_2	4	6	B_1, B_2, B_3
		2	C_3	3	8	A			
			C_{2v}	4	6	B_1, B_2			
			C_2	2	12	B			
			C_s	2	12	A', A''			
			T ₂	1	D_{2d}	8	3	B_2	
					C_{3v}	6	4	A_1	
		T _h	24	E _g	1	D_{2h}	8	3	A_g
						D_2	4	6	A
						C_i	2	12	A_g
					2	D_{2h}	8	3	A_u
						D_2	4	6	A
						C_i	2	12	A_u
T _g	1			D_{2h}	8	3	B_{1g}, B_{2g}, B_{3g}		
				S_6	6	4	A_g		
				D_2	3	8	B_1, B_2, B_3		
	2			C_3	3	8	A		
				C_i	2	12	A_g		
				T _u	1	D_{2h}	8	3	B_{1u}, B_{2u}, B_{3u}
S_6	6	4	A_u						
D_2	4	6	B_1, B_2, B_3						
2	C_3	3	8		A				
	C_i	2	12		A_u				
	O	24	E		1	D_4	8	3	A_1, B_1
2				D_2	4	6	A		
				C_2	2	12	A, B		

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term	
<i>I</i>	60	T_1	1	D_4	8	3	A_2	
				D_3	6	4	A_2	
			2	D_2	4	6	B_1, B_2, B_3	
				C_3	3	8	A	
				C_2	2	12	A, B	
				C_2	2	12	A, B	
		T_2	1	D_4	8	3	B_2	
				D_3	6	4	A_1	
			2	D_2	4	6	B_1, B_2, B_3	
				C_3	3	8	A	
				C_2	2	12	A, B	
				C_2	2	12	A, B	
		T_1	1	D_5	10	6	A_2	
				D_3	6	10	A_2	
			2	D_2	4	15	B_1, B_2, B_3	
				C_3	3	20	A	
				C_2	2	30	A, B	
				C_2	2	30	A, B	
		T_2	1	D_5	10	6	A_2	
				D_3	6	10	A_2	
			2	D_2	4	15	B_1, B_2, B_3	
				C_3	3	20	A	
				C_2	2	30	A, B	
				C_2	2	30	A, B	
G	1	T	12	5	A			
		D_3	6	10	A_1, A_2			
	2	D_2	4	15	B_1, B_2, B_3			
		C_3	3	20	A			
		C_2	2	30	A, B			
		C_2	2	30	A, B			
H	1	D_5	10	6	A_1			
		D_3	6	10	A_1			
	2	D_2	4	15	A, B_1, B_2, B_3			
		C_3	3	20	A			
		C_2	2	30	A, B			
		C_2	2	30	A, B			
T	12	E	1	D_2	4	3	A	
			1	D_2	4	3	B_1, B_2, B_3	
	D_{6h}	24	E_{1g}	1	D_{2h}	8	3	B_{2g}, B_{3g}
				2	D_2	4	6	B_2, B_3
D_{6h}	24	E_{1g}	2	C_{2v}	4	6	A_2, B_1, B_2	
				C_{2h}	4	6	A_g, B_g	
			3	C_2	2	12	A, B	
				C_i	2	12	A_g	
			E_{2g}	1	C_s	2	12	A', A''
					D_{2h}	8	3	A_g, B_{1g}
D_{6h}	24	E_{2g}	2	D_2	4	6	A, B_1	

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term		
<i>D</i> _{6d}	24	<i>E</i> _{1u}	3	<i>C</i> _{2v}	4	6	<i>A</i> ₁ , <i>A</i> ₂ , <i>B</i> ₂		
				<i>C</i> _{2h}	4	6	<i>A</i> _g , <i>B</i> _g		
				<i>C</i> ₂	2	12	<i>A</i> , <i>B</i>		
				<i>C</i> _i	2	12	<i>A</i> _g		
				<i>C</i> _s	2	12	<i>A'</i> , <i>A''</i>		
				<i>D</i> _{2h}	8	3	<i>B</i> _{2u} , <i>B</i> _{3u}		
			1	<i>D</i> ₂	4	6	<i>A</i> , <i>B</i> ₁		
				2	<i>C</i> _{2v}	4	6	<i>A</i> ₁ , <i>B</i> ₁ , <i>B</i> ₂	
					<i>C</i> _{2h}	4	6	<i>A</i> _u , <i>B</i> _u	
					<i>C</i> ₂	2	12	<i>A</i> , <i>B</i>	
				3	<i>C</i> _i	2	12	<i>A</i> _u	
					<i>C</i> _s	2	12	<i>A'</i> , <i>A''</i>	
			<i>E</i> _{2u}		1	<i>D</i> _{2h}	8	3	<i>A</i> _u , <i>B</i> _{1u}
				<i>D</i> ₂		4	6	<i>A</i> , <i>B</i> ₁	
				2	<i>C</i> _{2v}	4	6	<i>A</i> ₁ , <i>A</i> ₂ , <i>B</i> ₁	
					<i>C</i> _{2h}	4	6	<i>A</i> _u , <i>B</i> _u	
					3	<i>C</i> ₂	2	12	<i>A</i> , <i>B</i>
						<i>C</i> _i	2	12	<i>A</i> _u
		<i>E</i> ₁	2	<i>C</i> _s	2	12	<i>A'</i> , <i>A''</i>		
				<i>D</i> ₂	4	6	<i>B</i> ₂ , <i>B</i> ₃		
				<i>C</i> _{2v}	4	6	<i>B</i> ₁ , <i>B</i> ₂		
				3	<i>C</i> ₂	2	12	<i>A</i> , <i>B</i>	
					<i>C</i> _s	2	12	<i>A'</i> , <i>A''</i>	
					<i>D</i> _{2d}	8	3	<i>B</i> ₁ , <i>B</i> ₂	
			<i>E</i> ₂	1	<i>D</i> ₂	4	6	<i>A</i> , <i>B</i> ₁	
					<i>C</i> _{2v}	4	6	<i>A</i> ₁ , <i>A</i> ₂	
				2	<i>C</i> ₂	2	12	<i>A</i> , <i>B</i>	
					<i>C</i> _s	2	12	<i>A'</i> , <i>A''</i>	
				3	<i>D</i> ₆	12	2	<i>B</i> ₁ , <i>B</i> ₂	
					<i>C</i> _{6v}	12	2	<i>B</i> ₁ , <i>B</i> ₂	
			2		<i>D</i> ₂	4	6	<i>B</i> ₂ , <i>B</i> ₃	
				<i>C</i> _{2v}	4	6	<i>B</i> ₁ , <i>B</i> ₂		
				3	<i>C</i> ₂	2	12	<i>B</i>	
			<i>E</i> ₃	1	<i>D</i> _{2d}	8	3	<i>A</i> ₁ , <i>A</i> ₂	
					<i>D</i> ₂	4	6	<i>A</i> , <i>B</i> ₁	
				2	<i>C</i> _{2v}	4	6	<i>A</i> ₁ , <i>B</i> ₂	
		<i>C</i> ₂			2	12	<i>A</i> , <i>B</i>		
		3		<i>C</i> _s	2	12	<i>A'</i> , <i>A''</i>		
				<i>E</i> ₄	2	<i>D</i> ₂	4	6	<i>B</i> ₂ , <i>B</i> ₃
		<i>C</i> _{2v}	4			6	<i>B</i> ₁ , <i>B</i> ₂		
		3	<i>C</i> ₂		2	12	<i>A</i> , <i>B</i>		
		<i>E</i> ₅	2	<i>D</i> ₂	4	6	<i>B</i> ₂ , <i>B</i> ₃		
				<i>C</i> _{2v}	4	6	<i>B</i> ₁ , <i>B</i> ₂		

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term
			3	C_2	2	12	A, B
D_6	12	E_1	1	C_3	2	12	A', A''
			2	D_2	4	3	B_2, B_3
			2	C_2	2	6	A, B
D_{5h}	20	E_1'	1	D_2	4	3	A, B_1
			2	C_2	2	6	A, B
			2	C_2	2	6	A, B
			1	C_{2v}	4	5	A_1, B_1
			2	C_2	2	10	A, B
			2	C_3	2	10	A', A''
		E_2'	1	C_{2v}	4	5	A_1, B_1
			2	C_2	2	10	A, B
			2	C_3	2	10	A', A''
			1	C_{2v}	4	5	A_2, B_2
			2	C_2	2	10	A, B
			2	C_3	2	10	A', A''
D_{5d}	20	E_{1g}	1	C_{2h}	4	5	A_g, B_g
			2	C_2	2	10	A, B
			2	C_i	2	10	A_g
			2	C_3	2	10	A', A''
			1	C_{2h}	4	5	A_g, B_g
			2	C_2	2	10	A, B
		E_{2g}	1	C_{2h}	4	5	A_g, B_g
			2	C_2	2	10	A, B
			2	C_i	2	10	A_g
			2	C_3	2	10	A', A''
			1	C_{2h}	4	5	A_u, B_u
			2	C_2	2	10	A, B
E_{1u}	1	C_{2h}	4	5	A_u, B_u		
	2	C_2	2	10	A, B		
	2	C_i	2	10	A_u		
	2	C_3	2	10	A', A''		
	1	C_{2h}	4	5	A_u, B_u		
	2	C_2	2	10	A, B		
E_{2u}	1	C_{2h}	4	5	A_u, B_u		
	2	C_2	2	10	A, B		
	2	C_i	2	10	A_u		
	2	C_3	2	10	A', A''		
	1	C_{2h}	4	5	A_u, B_u		
	2	C_2	2	10	A, B		
D_5	10	E_1	1	C_2	2	5	A, B
			1	C_2	2	5	A, B
D_{4h}	16	E_g	1	D_{2h}	8	2	B_{2g}, B_{3g}
			2	D_2	4	4	B_2, B_3
			2	C_{2v}	4	4	B_1, B_2
		2	C_{2h}	4	4	B_g	
		3	C_2	2	8	B	
		1	D_{2h}	8	2	B_{2u}, B_{3u}	

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term
			2	D_2	4	4	B_2, B_3
				C_{2v}	4	4	B_1, B_2
				C_{2h}	4	4	B_u
			3	C_2	2	8	B
D_{4d}	16	E_1	2	D_2	4	4	B_2, B_3
				C_{2v}	4	4	B_1, B_2
			3	C_2	2	8	B
		E_2	1	D_4	8	2	B_1, B_2
				C_{4v}	8	2	B_1, B_2
			2	C_4	4	4	B
		E_3	2	D_2	4	4	B_2, B_3
				C_{2v}	4	4	B_1, B_2
			3	C_2	2	8	B
D_4	8	E	1	D_2	4	2	B_2, B_3
			2	C_2	2	4	B
D_{3h}	12	E'	1	C_{2v}	4	3	A_1, B_2
			2	C_2	2	6	A, B
				C_s	2	6	A', A''
		E''	1	C_{2v}	4	3	A_2, B_1
			2	C_2	2	6	A, B
				C_s	2	6	A', A''
D_{3d}	12	E_g	1	C_{2h}	4	3	A_g, B_g
			2	C_2	2	6	A, B
				C_i	2	6	A_g
				C_s	2	6	A', A''
		E_u	1	C_{2h}	4	3	A_u, B_u
			2	C_2	2	6	A, B
				C_i	2	6	A_u
				C_s	2	6	A', A''
D_3	6	E	1	C_2	2	3	A, B
D_{2d}	8	E	1	D_2	4	2	B_2, B_3
				C_{2v}	4	2	B_1, B_2
			2	C_2	2	4	B
S_{10}	10	E_{1g}	1	C_i	2	5	A_g
		E_{2g}	1	C_i	2	5	A_g
		E_{1u}	1	C_i	2	5	A_u
		E_{2u}	1	C_i	2	5	A_u
S_8	8	E_1	2	C_2	2	4	B
		E_2	1	C_4	4	2	B
		E_3	2	C_2	2	4	B
S_6	6	E_g	1	C_i	2	3	A_g
		E_u	1	C_i	2	3	A_u

• Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term
S_4	4	E	1	C_2	2	2	B
C_8	8	E_1	2	C_2	2	4	B
		E_2	1	C_4	4	2	B
C_{6h}	12	E_3	2	C_2	2	4	B
			1	C_{2h}	4	3	B_g
			2	C_2	2	6	B
				C_i	2	6	A_g
		E_{1g}		C_4	2	6	A''
			1	C_{2h}	4	3	A_g
			2	C_2	2	6	A
				C_i	2	6	A_g
		E_{2g}		C_4	2	6	A'
			1	C_{2h}	4	3	A_u
			2	C_2	2	6	A
				C_i	2	6	A_g
		E_{1u}		C_4	2	6	A'
			1	C_{2h}	4	3	B_u
2	C_2		2	6	B		
	C_i		2	6	A_u		
E_{2u}		C_4	2	6	A'		
	1	C_{2h}	4	3	A_u		
	2	C_2	2	6	A		
		C_i	2	6	A_u		
C_{6v}	12	E_1	1	C_6	2	6	A''
			2	C_2	2	6	A''
				C_3	2	6	A''
		E_2	1	C_{2v}	4	3	B_{1, B_2}
			2	C_2	2	6	B
				C_3	2	6	A', A''
				C_6	2	6	A', A''
		E_3	1	C_{2v}	4	3	A_1, A_2
			2	C_2	2	6	A
		C_6	6	E_1	1	C_2	2
E_2	1			C_2	2	3	A
E_3	1			C_3	2	5	A'
C_{5h}	10	E'_1	1	C_5	2	5	A'
		E'_2	1	C_5	2	5	A'
		E''_1	1	C_5	2	5	A''
		E''_2	1	C_5	2	5	A''
C_{5v}	10	E_1	1	C_5	2	5	A', A''
		E_2	1	C_5	2	5	A', A''
C_{4h}	8	E_g	1	C_{2h}	4	2	B_g
			2	C_2	2	4	B
		E_u	1	C_{2h}	4	2	B_u
			2	C_2	2	4	B
C_{4v}	8	E	1	C_{2v}	4	2	B_{1, B_2}
			2	C_2	2	4	B
C_4	4	E	1	C_2	2	2	B
C_{3h}	6	E'	1	C_3	2	3	A'

Table 1 (Continued)

Group	Order	Electron term	Level	Subgroup	Order of subgroup	Number of equivalent extremal points	Stable electron term
C_{3v}	6	E''	1	C_s	2	3	A''
		E	1	C_s	2	3	A', A''

a) Notation for irreducible representations:

1. A : 1-Dimensional representation which is symmetric with respect to rotation about the principal axis;
 B : 1-dimensional representation which is antisymmetric with respect to rotation about the principal axis;
 E : 2-dimensional representation;
 T : 3-dimensional representation;
 G : 4-dimensional representation;
 H : 5-dimensional representation.
 2. Subscript 1, 2:
 For A and B representations only-symmetric (1) or antisymmetric (2) with respect to C_2 axis perpendicular to the principal axis, or, in the absence of this, with respect to a vertical symmetry plane.
 3. If there exist higher subscripts (3, 4 or 5), these indicate the different symmetry with respect to specific symmetry operation, which depends on the given symmetry group (for example in the groups D_{6d} , D_{4d} , D_{2h} , S_8).
 4. Primes and double primes (' and ''):
 Where appropriate, these indicate symmetry (') or antisymmetry ('') with respect to a horizontal mirror plane.
 5. g and u subscripts:
 Where appropriate, these denote symmetry (g — gerade) or antisymmetry (u — ungerade) with respect to an inversion centre.
- If these rules allow several different labels, g and u take precedence over 1 and 2, which take precedence over ' and ''

Discussion

We start our discussion by presenting the scheme for the octahedral symmetry group (O_h). Its order is equal to 48 and contains 6 multi-dimensional irreducible representations (E_g , E_u , T_{1g} , T_{2g} , T_{1u} , and T_{2u}). It contains all together 24 subgroups (O , T_d , T_h , D_{4h} , D_{3d} , D_4 , D_3 , D_2 , D_{2h} , D_{2d} , S_6 , S_4 , T , C_{3v} , C_{2h} , C_{4v} , C_{4h} , C_{2v} , C_4 , C_3 , C_2 , C_i , C_s , and C_1). The following subgroups fulfil the three above-mentioned conditions: D_{4h} , D_{3d} , D_4 , D_3 , D_2 , D_{2h} , D_{2d} , S_6 , C_3 , C_2 , C_{2h} , C_{2v} , C_i , and C_s . In the case of three-times degenerate T_{2g} electron term this is in the first level split to B_{2g} and E_g terms in the subgroup D_{4h} or to terms A_{1g} and E_g in the subgroup D_{3d} . The

other three direct subgroups (O , T_h , and T_d) do not cause splitting of T_{2g} electron term. In the second level the degenerate representations of the first level split to the following nondegenerate irreducible representations: $B_2(D_4)$, $A_1(D_3)$, $B_1(D_{2d})$, B_{1g} , B_{2g} , and $B_{3g}(D_{2h})$, $A_g(S_6)$, $A_1(C_{3v})$, A_g and $B_g(C_{2h})$. In the third level the degenerate irreducible representations of the second level split to the following nondegenerate irreducible representations: B_1 , B_2 , and $B_3(D_2)$, $A(C_3)$, B_1 and $B_2(C_{2v})$, $B_g(C_{2h})$, A and $B(C_2)$, $A_g(C_i)$, A' and $A''(C_s)$. In the fourth level we have only one nondegenerate irreducible representation $B(C_2)$.

If we limit ourselves to the so-called Jahn—Teller active vibrations which give already the nonzero contribution to vibronic interaction in the first order of perturbation theory, it is sufficient to take into account only the centro-symmetric vibrations. In this case we can limit ourselves only to the subgroups which possess the inversion symmetry operation. Such are only the subgroups D_{4h} , D_{3d} , S_6 , D_{2h} , C_{2h} , and C_i . Through this limitation we get in the first level the irreducible representations $B_{2g}(D_{4h})$ and $A_{1g}(D_{3d})$. In the second level we get B_{1g} , B_{2g} , and $B_{3g}(D_{2h})$, A_g and $B_g(C_{2h})$, and $A_g(S_6)$ irreducible representations. In the third level we have only $B_g(C_{2h})$ and $A_g(C_i)$ irreducible representations. For octahedral system in three-times degenerate T_{2g} electron term using this limitation we get:

i) three equivalent geometries of D_{4h} symmetry	}	1st level
ii) four equivalent geometries of D_{3d} symmetry		
iii) six equivalent geometries of D_{2h} symmetry		
iv) eight equivalent geometries of S_6 symmetry	}	2nd level
v) twelve equivalent geometries of C_{2h} symmetry		
vi) twelve equivalent geometries of C_{2h} symmetry	}	3rd level
vii) twenty-four equivalent geometries of C_i symmetry		

It is necessary to note that conclusions based on assuming only the centro-symmetric subgroups are valid only in the first order of perturbation theory. Therefore in some approaches [1—19] this limitation occurs.

The classification of molecular symmetry point groups according to given rules (presented in Table 1) determines all the possible symmetries of systems which undergo the Jahn—Teller deformation. The presented table therefore represents the complete classification of possible symmetries of Jahn—Teller systems for all symmetry point groups of molecules. Thus the existence of deformations of formerly symmetric system into the system of different symmetry cannot be due to the Jahn—Teller effect. On the basis of group theory we can reject also some commonly used explanations of the transition between some geometries as the consequence of Jahn—Teller effect. Therefore for example the transition tetrahedron (symmetry T_d) \leftrightarrow square planar (symmetry D_{4h}) cannot be due to the Jahn—Teller effect, because the group D_{4h} (order 16) is not subgroup of T_d group (order 24). Besides, in some recent papers [19] this transition is explained as due to the Jahn—Teller effect. Similarly the transition of five-coordinated complexes

between trigonal bipyramid (symmetry D_{3h}) and square pyramid (symmetry C_{4v}) is not due to Jahn—Teller deformation because group C_{4v} (order 8) is not subgroup of D_{3h} group (order 12).

On the other hand, we must have in mind also the limitation of group-theoretical approach to the solution of problems of symmetries of Jahn—Teller systems. Such an approach on one side gives the geometries of all possible extremal points of adiabatic potential for given system but on the other side it says nothing whether these extremal points will be realized or not.

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